# Linear algebra problems arising in discontinuous Galerkin finite element discretizations 

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## Outline

- Introduction to discontinuous Galerkin (DG) FEM
- DG-FEM for the Poisson problem
- formulation of DG methods
- matrix form
- survey of the literature on solvers
- DG-FEM for Maxwell's problems
- mixed, indefinite and eigenvalue problems


## Introduction to DG-FEM

DG-FEM are finite element methods based on completely discontinuous finite element spaces

Ingredients

- $\mathcal{T}_{h}=\{K\}$ partition of the domain $\Omega$
- $\mathcal{P}^{\ell}\left(\mathcal{T}_{h}\right)=$ piecewise polynomials of degree $\ell$ in each element possibly discontinuous across interelement boundaries
- Local variational formulation (element-by-element) $\rightarrow$ interelement continuity conditions imposed within the variational formulation (no special boundary degrees of freedom, no Lagrange multipliers)


## Introduction to DG-FEM

## Main Features

- Wide range of PDE's treated within the same unified framework
- Flexibility in the mesh design $\rightarrow$ good for adaptivity
- non-matching grids (hanging nodes)
- non-uniform approximation degrees

- Block-diagonal (even diagonal) mass matrices
- Drawback: high number of degrees of freedom


## DG for the Poisson Problem

## The Poisson Problem

Given $f \in L^{2}(\Omega)$ and $g_{\mathcal{D}} \in H^{1 / 2}(\Omega)$, find $u \in H^{1}(\Omega)$ s.t.

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\begin{aligned}
-\Delta u & =f & & \text { in } \Omega \subset \mathbb{R}^{2} \\
u & =g_{\mathcal{D}} & & \text { on } \partial \Omega
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## Variational Formulation

Find $u \in H^{1}(\Omega)$ with $u=g_{\mathcal{D}}$ on $\partial \Omega$ s.t.

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\int_{\Omega} \nabla u \cdot \nabla v=\int_{\Omega} f v \quad \forall v \in H_{0}^{1}(\Omega)
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For continuous FE spaces: $\int_{\Omega} \nabla u_{h} \cdot \nabla v_{h}=\int_{\Omega} f v_{h} \rightarrow$ standard FEM

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- $\mathcal{T}_{h}$ partition of the domain; $\mathcal{F}_{h}$ set of all edges (faces in 3D); $\mathcal{F}_{h}^{\mathcal{I}}, \mathcal{F}_{h}^{\mathcal{B}}$ sets of all interior and boundary edges, resp.


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- Multiply by $v_{h} \in V_{h}$ and integrate by parts in each element

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- Sum over all elements

$$
\sum_{K \in \mathcal{T}_{h}} \int_{K} \nabla u \cdot \nabla v_{h}-\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \nabla u \cdot \mathbf{n}_{K} v_{h}=\int_{\Omega} f v_{h}
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\sum_{K \in \mathcal{I}_{h}} \int_{K} \nabla u \cdot \nabla v_{h}-\sum_{K \in \mathcal{I}_{h}} \int_{\partial K} \nabla u \cdot \mathbf{n}_{K} v_{h}=\int_{\Omega} f v_{h}
$$

Key formula

$$
\sum_{K \in \mathcal{T}_{h}} \int_{\partial K} \nabla u \cdot \mathbf{n}_{K} v_{h}=\sum_{f \in \mathcal{F}_{h}} \int_{f}\{\nabla u\} \cdot \llbracket v_{h} \rrbracket_{N}+\underbrace{\sum_{f \in \mathcal{F}_{h}^{\mathcal{I}}} \int_{f} \llbracket \nabla u \rrbracket_{N}\left\{\left\{v_{h}\right\}\right\}}_{=0}
$$

Averages and jumps on interior edges

- $\{v\}\}:=\left(v^{+}+v^{-}\right) / 2$
- $\llbracket v \rrbracket_{N}:=v^{+} \mathbf{n}^{+}+v^{-} \mathbf{n}^{-}$

$$
\begin{aligned}
\{\mathbf{q}\} & :=\left(\mathbf{q}^{+}+\mathbf{q}^{-}\right) / 2 \\
\llbracket \mathbf{q} \rrbracket_{N} & :=\mathbf{q}^{+} \cdot \mathbf{n}^{+}+\mathbf{q}^{-} \cdot \mathbf{n}^{-}
\end{aligned}
$$

Averages and jumps on boundary edges

- $\{\mathbf{q}\}=\mathbf{q}$
- $\llbracket v \rrbracket_{N}:=v \mathbf{n}, \llbracket u \rrbracket_{N}:=\left(u-g_{\mathcal{D}}\right) \mathbf{n}$



## DG for the Poisson Problem

DG Methods of the Interior Penalty Family (IP-DG)

$$
\sum_{K \in \mathcal{T}_{h}} \int_{K} \nabla u_{h} \cdot \nabla v_{h}-\sum_{f \in \mathcal{F}_{h}} \int_{f}\left\{\left\{\nabla u_{h}\right\}\right\} \cdot \llbracket v_{h} \rrbracket_{N}
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& -k \sum_{f \in \mathcal{F}_{h}} \int_{f} \llbracket u_{h} \rrbracket_{N} \cdot\left\{\left\{\nabla_{h} v_{h}\right\}\right. & =\int_{\Omega} f v_{h}
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## Matrix Form

IP-DG Methods
$\sum_{K \in \mathcal{T}_{h}} \int_{K} \nabla u_{h} \cdot \nabla v_{h}$


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- $A$ large, sparse, $\kappa_{2}(A) \sim h^{-2}$


## Matrix Form



## Matrix Form

$$
V+\alpha S
$$



## Matrix Form

$$
V-F-k F^{T}+\alpha S=: A
$$



## Solvers

- Domain decomposition (Schwarz) preconditioners
[Rusten, Vassilevski \& Winther, 1996], [Feng \& Karakashian, 2001],
[Lasser \& Toselli, 2003], [Brenner \& Wang, 2005],
[Antonietti \& Ayuso, 2007-08], [Dryja, Galvis \& Sarkis, 2007]
- Multigrid methods
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- Multilevel preconditioners with cG discretization at lowest level [Warsa, Benzi, Wareing \& Morel, 2004], [Antonietti \& Ayuso, 2007]
- Norm preconditioners
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[Georgoulis \& Loghin, 2008]: norm matrix as precond.


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NIP-DG provides an example in which, although the symmetric part of the operator has negative eigenvalues, the convergence of GMRES takes place all the same.

## Maxwell's Problems

- Mixed problem (low-frequency approximation of the time-harmonic Maxwell's problem in insulating materials)

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\begin{array}{ll}
\nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)-\varepsilon \nabla p=\mathbf{f} & \text { in } \Omega \\
\nabla \cdot(\varepsilon \mathbf{u})=0 & \text { in } \Omega \\
\mathbf{n} \times \mathbf{u}=\mathbf{0}, \quad p=0 & \text { on } \partial \Omega
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## Maxwell's Problems

## DG Discretizations

- Key ingredient: DG discretization of the curl-curl operator
- similar to DG discretization of the Laplacian (here: penalization of the tangential jumps of $\mathbf{u}$ )
- complex vector-valued fields
- "large" kernel


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- imposition of the divergence-free constraint in a DG fashion
- penalization of the jumps of $p$
- balancing $\mathbf{V}_{h}$ and $Q_{h}$ (polynomials of degree $\ell$ for $\mathbf{u}$ and $\ell+1$ for $p$ )
- theoretical analysis based on an underlying stable discretization with conforming elements [Houston, Perugia \& Schötzau, 2005]


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\nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)-\omega^{2} \varepsilon \mathbf{u}=\mathbf{f}
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- DG discretization straightforward (recall: block diagonal mass matrix)
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## Matrix Form

Mixed Problem

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\left[\begin{array}{cc}
A & B^{T} \\
B & -\beta C
\end{array}\right]\left[\begin{array}{l}
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\mathbf{p}
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- $A$ and $C$ are positive semidefinite ("large" kernel)
- Non-singular overall system


## Matrix Form

Indefinite Problem and Eigenvalue Problem

$$
\nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)-\omega^{2} \varepsilon \mathbf{u}=\mathbf{f} \quad \nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)=\omega^{2} \varepsilon \mathbf{u}
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$$
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\nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)-\omega^{2} \varepsilon \mathbf{u}=\mathbf{f} & \nabla \times\left(\mu^{-1} \nabla \times \mathbf{u}\right)=\omega^{2} \varepsilon \mathbf{u} \\
{\left[A-\omega^{2} M\right] \mathbf{u}=\mathbf{f}} & A \mathbf{u}=\omega^{2} M \mathbf{u}
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- $A-\omega^{2} M$ non singular (indefinite), provided that $\omega^{2}$ is not a discrete eigenvalue of the pencil $(A \mid M)$


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Literature on conforming methods (almost nothing available for DG)

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- preconditioned shift-and-invert Lanczos, Jacobi-Davidson [Arbenz \& Geus, 1999-2005], [Simoncini, 2003]


## Conclusions

- For the Poisson problem: extension to DG-FEM of solution techniques already studied for conforming methods with variants either in their formulation or in the analysis


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-\Delta u-k^{2} u=0 \quad \text { in } \Omega, \quad \nabla u \cdot \mathbf{n}-i \omega u=g \quad \text { on } \partial \Omega
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with local approximating spaces made of linear combinations of plane waves $\sum_{k=1}^{p} a_{k} \exp \left(i \omega \mathbf{d}_{k} \cdot \mathbf{x}\right)$ (instead of polynomials); here, the choice of the plane wave directions $\mathbf{d}_{k}$ also affects the conditioning

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- DG-FEM provide "complicated" test cases

